Cosmic Rays During the Recent Unusual Solar minimum

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Observation

Intensity of Galactic cosmic rays (GCRs) measured during the recent solar minimum was the highest ever recorded. We studied CRs data on spacecrafts near the Earth and ground-based neutron monitors, which indicate that the modulation of CRs is not dominated by the mechanism of particle drift through current sheet during this A < 0 cycle as we normally think.

Modulation

We use a model of GCRs transport in the three-dimensional heliosphere based on a simulation of Markov stochastic process to study the possible reasons. Our preliminary results show that it is due to the weaker Interplanetary magnetic field and lower perpendicular diffusion coefficient.





• Solar magnetic field



• *A*⁺ are times when the solar magnetic field is directed outward from the sun in the northern polar and inward in the southern polar region.





Transport through a wavy current sheet



 Diffusion with respect to the Parker spiral (left). The global drift pattern of positively charged particles in A > 0 and A < 0 solar magnetic epoch (right).



Solar modulation of GCRs



• Both helium and electrons vary in anti-correlation with 11-year solar activity cycle. In the A < 0, the time profiles of positively charged particles peaked, whereas they were more or less flat in A > 0.

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Figure 1: monthly mean rates for NMs from 1980.01 to 2011.01; rates are normalized to 100% for February of 1987.





- Rome 6.27Gv
- Apatity 5.6Gv
- Hermanus 4.58Gv
- Moscow 2.43Gv
- Oulu 0.8Gv

- Jungfraujoch 4.49Gv
- Climax 3Gv
- Kiel 2.36Gv
- Magadan 2.09Gv



Figure 2: rates are normalized to 100% for March of 1987

Neavy ions on the ACE





- The number is the proton number of partilce, they are C,N,O,Ne,Na,Mg,Si,S,Fe.
- I use annual mean hourly rates data, this is the solar minimum of cycle 24(2009).
- Each ion has seven energy ranges. We calculation annual deviation to study.



Figure 3: We calculate deviation every year for all heavy ions, $\sqrt{\frac{(C-\bar{C})^2}{N}}$. The deviation vary in anti-correlation with 11-year solar activity cycle. And the deviation in 24 solar cycle is much bigger and sharp than the one in 23 solar cycle.



Figure 4: top:action; bottom:deviation compare

- We calculate annual deviation for all particles (5-28) from 1997 to 2010(blue line in bottom panel)
- We use f' = T(z)f to concentrate flux of all particles together.

$$T(z) = 10^{k(z-z_0)}$$
 (1)

• Deviation after processing is the red line in bottom panel. It is much smaller comparing with before.



Figure 5: flux before processing in the solar minimum 23

Figure 6: flux after processing





Figure 7: We calculate relative deviation between our modulated results and the observation for all elements. $D = \sqrt{\frac{\sum_{year} (\frac{fm-f_0}{f_0})^2}{years}}$. Except iron and some rare ions our modulation fit observation well.







Figure 8: Three Ulysses fast latitude scans:first and third take place at solar minimum and second one under solar maximum.



Figure 9: From top to bottom are shown the daily averaged flux of 38-125 Mev protons from 1990 to 2009 and the monthly averaged after data processing.

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 Monthly averaged data of 400Mev proton on the IMP-8 after dealing.

 The energy of Proton on the IMP-8 is from 70Mev to 400Mev, time is from 1980 to 2001. proton--398.200 Mev



Figure 10: Monthly averaged after dealing



Figure 11: Daily averaged data of 176Mev proton on voyager1, available energy is from 22Mev to 176Mev.

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Figure 12: blue:radius, red:latitude, green:longitude

Sokker-Planck transport equation

Fokker-Planck equation

$$\frac{\partial f}{\partial t} = -(V + V_d) \cdot \nabla f + \nabla \cdot (\kappa \cdot \nabla f) + \frac{p}{3} (\nabla \cdot V) \frac{\partial f}{\partial p}$$
(2)

with

Orift speed:

$$V_{dr} = \frac{V_{\rho}\rho}{3q} \nabla \times \frac{B}{B^2}$$
(3)

Oiffusion coefficient:

$$\kappa_{\parallel} = \kappa_{\parallel 0} \beta \left(\frac{p}{p_0}\right)^{b_{\parallel}} \left(\frac{B_e}{B}\right)^{a_{\parallel}}$$
(4)
$$\kappa_{\perp} = \kappa_{\perp 0} \beta \left(\frac{p}{p_0}\right)^{b_{\perp}} \left(\frac{B_e}{B}\right)^{a_{\perp}}$$
(5)







The current sheet

$$\theta_0 = \frac{\pi}{2} + \alpha \sin\left(\phi + \frac{r\Omega}{V_w}\right) \tag{6}$$

calculation of IMF

$$B = \frac{A}{r^2} (\hat{\boldsymbol{e}}_r - \Gamma \boldsymbol{e}_{\phi}) \left[1 - 2H(\theta - \theta_0) \right]$$
(7)

drift of wavy CS

$$v_{cs} = \frac{v}{6} \cos \alpha \frac{\triangle \theta_{cs}}{\sin(\alpha + \triangle \theta_{cs})}$$
(8)

with

$$\triangle \theta_{cs} = P/(80 \cos \alpha) \tag{9}$$

Stochastic differential equation



We use Markov stochastic process to solve the transport equation, X(s), Y(s), Z(s), P(s) are time-backward stochastic processes described by ($U = V + V_d$)

$$d\mathbf{x} = \left(\frac{\partial \kappa_{\perp}}{\partial \mathbf{x}} - U_{\mathbf{x}}\right) \, \mathbf{ds} + \sqrt{2\kappa_{\perp}} \, \mathbf{dW}_{\mathbf{x}}(\mathbf{s}) \tag{10}$$

$$dy = \left(\frac{\partial \kappa_{\perp}}{\partial y} - U_{y}\right) ds + \sqrt{2\kappa_{\perp}} dW_{y}(s)$$
(11)

$$dz = \left(\frac{\partial \kappa_{\parallel}}{\partial z} - U_z\right) ds + \sqrt{2\kappa_{\parallel}} dW_z(s)$$
(12)

$$\mathrm{d}\boldsymbol{p} = \frac{\boldsymbol{p}}{3} \left(\nabla \cdot \boldsymbol{V} \right) \, \mathrm{d}\boldsymbol{s} \tag{13}$$

$$f(x, y, z, p, t) = \left\langle \int_0^t Q(X(s), Y(s), Z(s), P(s), s) \, \mathrm{d}s \right\rangle$$
(14)

Sample trajectory for *qA* < 0





- A simulated stochastic process with Magnetic field polarity qA < 0.
- The radial distance, latitude, and momentum of the simulated particle are shown as functions of backward time S.
- The process starts at 1AU in the equatorial with 1GeVc⁻¹ momentum and runs backward in time until it exits at 75 AU.

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Latitude-longitude distribution



0.6 qA=-1 0.4 0.2 latitude 0 -0.2 -0.4 -0.6 -3 -2 2 3 -1 0 -4 longitude

longitude and latitude distribution

Figure 14: Latitude-longitude distribution of test particles for qA < 0.

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- The same as figure 1, but for qA > 0
- We note that the particle gains more momentum than in the case of qA negative.

S Latitude-longitude distribution





Figure 15: Latitude-longitude distribution of test particles for qA < 0.

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Figure 16: Ideal model

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Cosmic Ray Modulation

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Figure 17: We use transport parameters carrying out a half year average

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- We model proton intensity according to Ulysses trace
- Observation circles are obtained from Ulysses COSPIN.







- We use transport parameters carrying out a half year average of each interval in our model of each solar minima.
- The observation data are obtained from neutron monitor and IMP-8.