

# Basic Space Plasmas Physics

## Assignment 4

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(Solve **three** out of the four problems! If time permits, you can solve all.)

### 1 Proof

1.

$$\begin{aligned}
 & -\frac{1}{2}v^2 [\nabla \cdot (\rho\vec{v})] - \rho\vec{v} \cdot [(\vec{v} \cdot \nabla)\vec{v}] + \nabla \cdot (\frac{1}{2}\rho v^2 \vec{v}) \\
 & = \rho\vec{v} \cdot \nabla(\frac{1}{2}v^2) - \rho\vec{v} \cdot [(\vec{v} \cdot \nabla)\vec{v}] \\
 & = \rho\vec{v} \cdot [\vec{v} \times (\nabla \times \vec{v})] \\
 & = 0
 \end{aligned} \tag{1.1}$$

2.

$$\begin{aligned}
 & \vec{v} \cdot (\nabla \cdot \Pi) - \nabla \cdot [\vec{v} \cdot \Pi] \\
 & = v_i \frac{\partial}{\partial x_j} \Pi_{ji} - \frac{\partial}{\partial x_j} (v_i \Pi_{ij}) \\
 & = v_i \frac{\partial}{\partial x_j} \Pi_{ji} - v_i \frac{\partial}{\partial x_j} \Pi_{ij} - \Pi_{ij} \frac{\partial}{\partial x_j} v_i \\
 & = -\Pi_{ij} \frac{\partial}{\partial x_j} v_i \quad (\Pi \text{ is a symmetric matrix}) \\
 & = -\nabla\vec{v} : \Pi \quad (\text{or } -\Pi : \nabla\vec{v})
 \end{aligned} \tag{1.2}$$

3.

$$\begin{aligned}
 & \vec{H} \cdot [\nabla \times (\vec{v} \times \vec{H})] + \vec{v} \cdot [(\nabla \times \vec{H}) \times \vec{H}] + \nabla \cdot [\vec{H} \times (\vec{v} \times \vec{H})] \\
 & = \vec{H} \cdot [\nabla \times (\vec{v} \times \vec{H})] + \vec{v} \cdot [(\nabla \times \vec{H}) \times \vec{H}] + (\nabla \times \vec{H}) \cdot (\vec{v} \times \vec{H}) - \vec{H} \cdot [\nabla \times (\vec{v} \times \vec{H})] \\
 & = 0
 \end{aligned} \tag{1.3}$$

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4.

$$\begin{aligned}
& \vec{H} \cdot \nabla^2 \vec{H} - \nabla \cdot [\vec{H} \times (\nabla \times \vec{H})] \\
&= \vec{H} \cdot \nabla^2 \vec{H} - (\nabla \times \vec{H}) \cdot (\nabla \times \vec{H}) + \vec{H} \cdot [\nabla \times (\nabla \times \vec{H})] \\
&= \vec{H} \cdot \nabla^2 \vec{H} - |\nabla \times \vec{H}|^2 + \vec{H} \cdot [\nabla(\nabla \cdot \vec{H}) - \nabla^2 \vec{H}] \\
&= -|\nabla \times \vec{H}|^2
\end{aligned} \tag{1.4}$$

## 2 Magnetohydrodynamics Energy Equation

The continuity equation is given by

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \tag{2.1}$$

where  $\rho$  is the mass density and  $\vec{v}$  is the flow velocity.

The momentum equation (equation of motion) is given by

$$\rho \left( \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \right) = \frac{\mu}{4\pi} (\nabla \times \vec{H}) \times \vec{H} - \nabla p + \nabla \cdot \Pi \tag{2.2}$$

where  $\vec{H}$  is magnetic field intensity,  $\mu$  is a material dependent parameter called the permeability,  $p$  is the scalar pressure, and  $\Pi$  is the viscous stress tensor.

$$\Pi_{ik} = \zeta \left( \frac{\partial v_i}{\partial x_k} + \frac{\partial v_k}{\partial x_i} - \frac{2}{3} \delta_{ik} \nabla \cdot \vec{v} \right) + \zeta' \delta_{ik} \nabla \cdot \vec{v} \tag{2.3}$$

The energy equation (conservation of energy) is given by

$$\frac{\partial W}{\partial t} + \nabla \cdot \vec{q} = 0 \tag{2.4}$$

where  $W$  is total energy per unit volume and  $\vec{q}$  is the energy flux density through the boundary of the fluid element. The total energy is the sum of the kinetic, magnetic, and internal energies:

$$W = \frac{1}{2} \rho v^2 + \frac{\mu}{8\pi} H^2 + \rho e \tag{2.5}$$

where  $e$  is the internal energy per unit mass. The energy flux density is given by

$$\begin{aligned}
\vec{q} &= \rho \vec{v} \left( \frac{1}{2} v^2 + e + \frac{p}{\rho} \right) + \frac{\mu}{4\pi} \vec{H} \times (\vec{v} \times \vec{H}) \\
&\quad - \frac{1}{\sigma} \left( \frac{c}{4\pi} \right)^2 \vec{H} \times (\nabla \times \vec{H}) - \vec{v} \cdot \Pi - \chi \nabla T
\end{aligned} \tag{2.6}$$

where  $\sigma$  is the electrical conductivity,  $\chi$  is the thermal conductivity, and  $T$  is the temperature.

Derive the non-conservation form of the energy equation as the following, which only contains the time derivative of  $e$ .

$$\rho \frac{\partial e}{\partial t} + \rho \vec{v} \cdot \nabla e = -p \nabla \cdot \vec{v} + \nabla \vec{v} : \Pi + \nabla \cdot (\chi \nabla T) + \frac{1}{\sigma} J^2 \tag{2.7}$$

**Hint:** The magnetohydrodynamics equations are all shown in Gaussian electromagnetic units system. The equations of magnetic field are given by

$$\frac{\partial \vec{H}}{\partial t} = \nabla \times (\vec{v} \times \vec{H}) + \frac{c^2}{4\pi\mu\sigma} \nabla^2 \vec{H} \tag{2.8}$$

$$\nabla \cdot \vec{H} = 0 \tag{2.9}$$

### 3 Frozen Flux Theorem

Under ideal MHD assumption, the magnetic field equations (SI units) reduce to

$$\frac{\partial \vec{B}}{\partial t} = \nabla \times (\vec{v} \times \vec{B}) \quad (3.1)$$

$$\nabla \cdot \vec{B} = 0 \quad (3.2)$$

where  $\vec{B}$  is the magnetic flux density. Consider a closed curve  $C$  within the fluid, and let every point on the curve be moving with the local fluid velocity. We say that  $C$  is co-moving with the fluid, in the Lagrangian sense. Let  $S$  be a surface boundary by  $C$ , then we can obtain the surface integral:

$$\iint \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} - \oint (\vec{v} \times \vec{B}) \cdot d\vec{C} = 0 \quad (3.3)$$

As  $S_1(t)$  moves in a infinite short time  $dt$ ,  $S_1(t)$ ,  $S_2(t + dt)$ , and swept lateral area  $S_3$  form a closed surface. This is shown in Figure 1 at time  $t + dt$ . For  $\nabla \cdot \vec{B} = 0$ , we can obtain that

$$\iint_{S_2} \vec{B}(t + dt) \cdot d\vec{S} - \iint_{S_1} \vec{B}(t + dt) \cdot d\vec{S} + \iint_{S_3} \vec{B}(t + dt) \cdot d\vec{S} = 0 \quad (3.4)$$

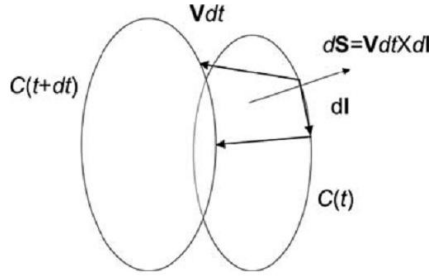


Figure 1: A surface moving with the fluid

Let  $dt \rightarrow 0$ , then

$$\frac{d}{dt} \iint_S \vec{B} \cdot d\vec{S} - \iint_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} + \oint_C (\vec{v} \times \vec{B}) \cdot d\vec{C} = 0 \quad (3.5)$$

With (3.3), we find

$$\frac{d}{dt} \iint_S \vec{B} \cdot d\vec{S} = 0 \quad (3.6)$$

We conclude that in ideal MHD, the magnetic flux through any co-moving closed circuit remains constant. This important result is called the frozen flux condition.

Consummate all the derivation.

### 4 Magnetic Diffusion

A unidirectional magnetic field  $\vec{B} = B(x, t)\vec{e}_y$  has the initial form

$$B(x, 0) = \begin{cases} +B_0, & x > 0 \\ -B_0, & x < 0 \end{cases} \quad (4.1)$$

If the magnetic Reynolds number is very small, according to the induction equation, the evolution of the magnetic field can be described as

$$\frac{\partial B}{\partial t} = \eta \frac{\partial^2 B}{\partial x^2} \quad (4.2)$$

where  $\eta$  is the magnetic viscous coefficient. The solution is

$$B(x, t) = B_0 \operatorname{erf}(\xi) \quad (4.3)$$

where  $\xi = \frac{x}{\sqrt{4\eta t}}$  and  $\operatorname{erf}(\xi)$  is the error function given by

$$\operatorname{erf}(\xi) = \frac{2}{\sqrt{\pi}} \int_0^\xi e^{-z^2} dz \quad (4.4)$$

Having known  $B(2\sqrt{4\eta t}, t) \approx 0.995B_0$  and  $\int_{-2}^2 [1 - \operatorname{erf}^2(\xi)] d\xi \approx 1.592$ , estimate the dissipation rate of the magnetic energy.